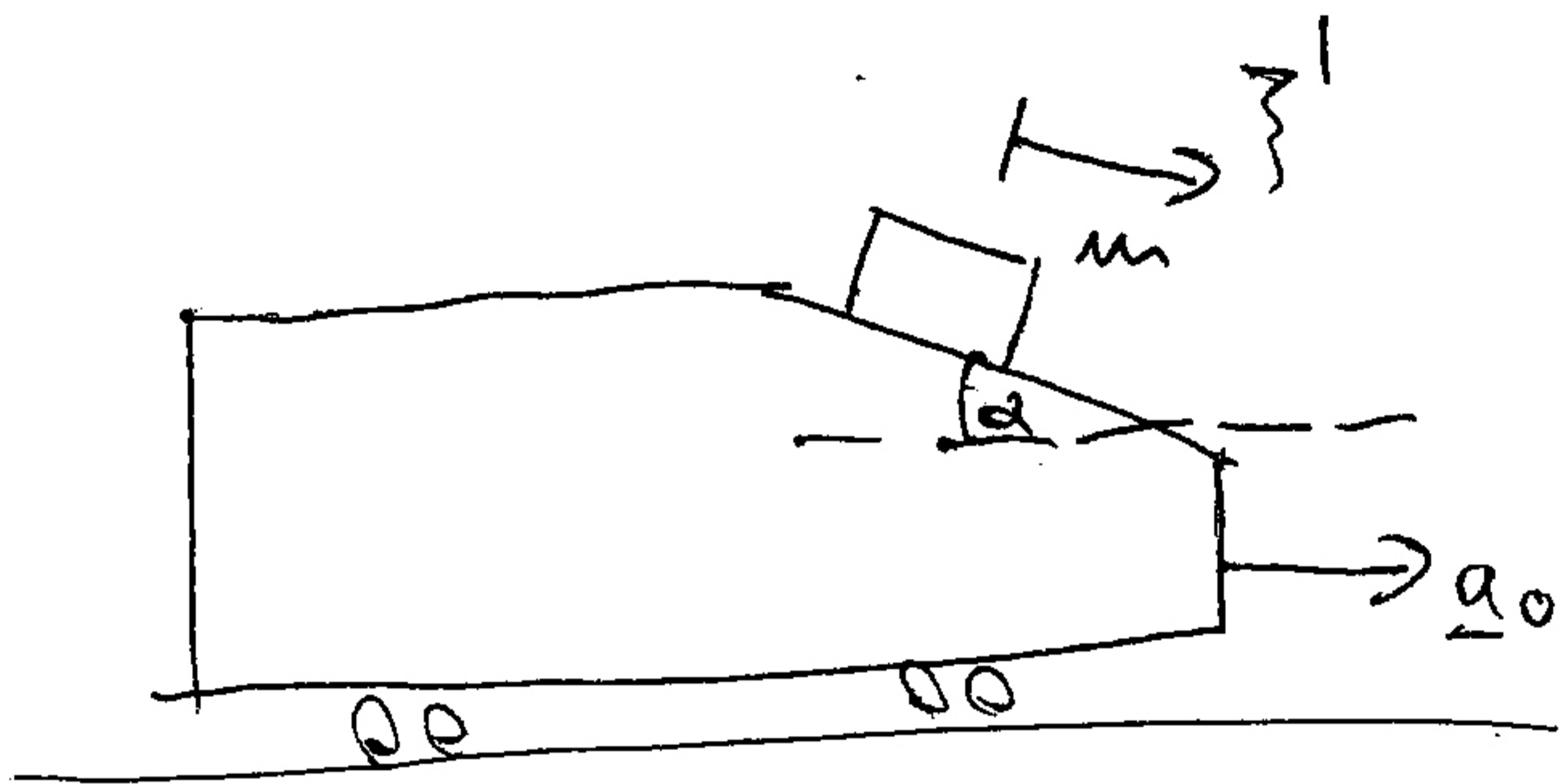


1. Legen a test gysvulsa a locitor viorvlytra 0:

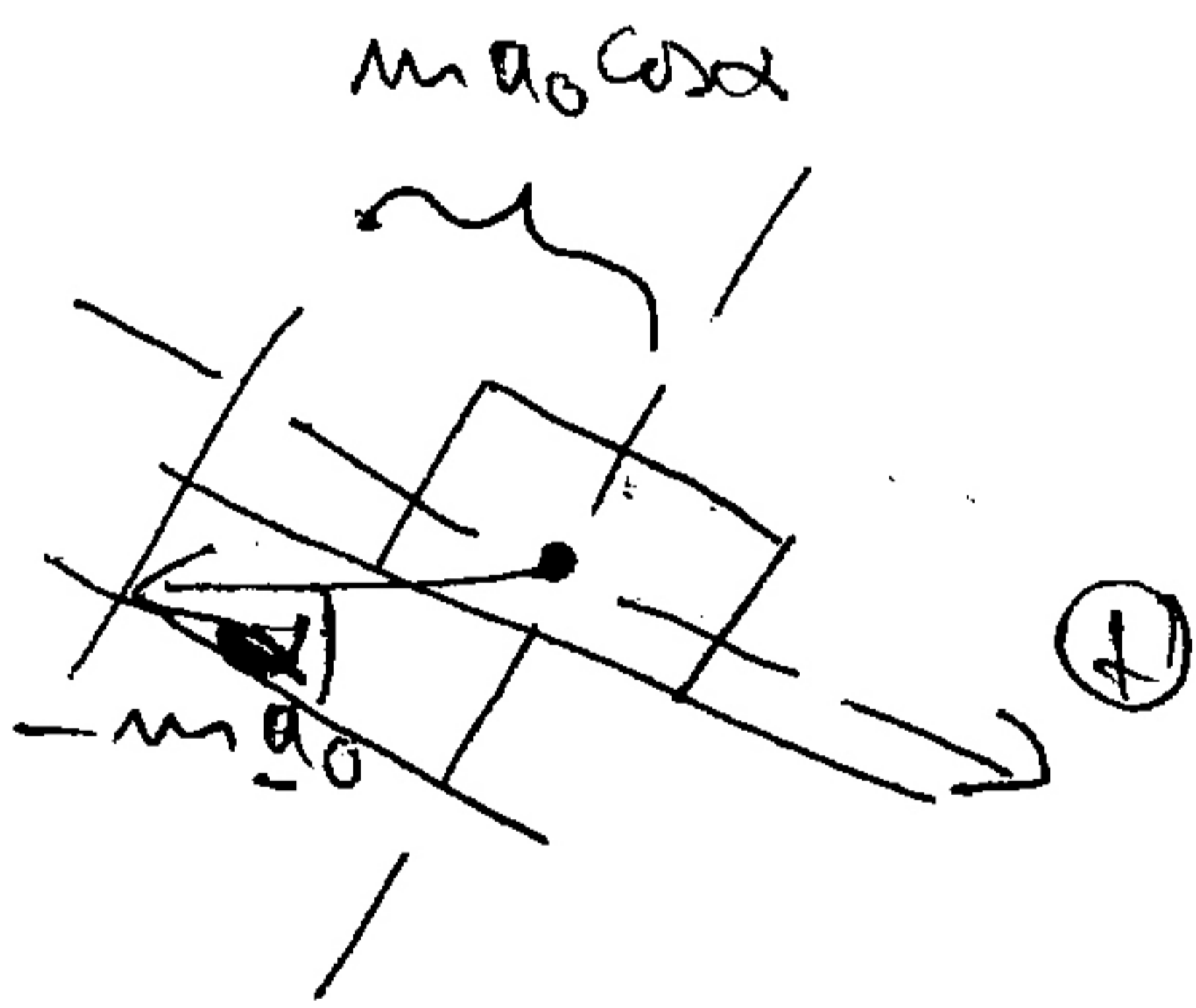
ellor, sddeberig hlydr,
neu modal el.



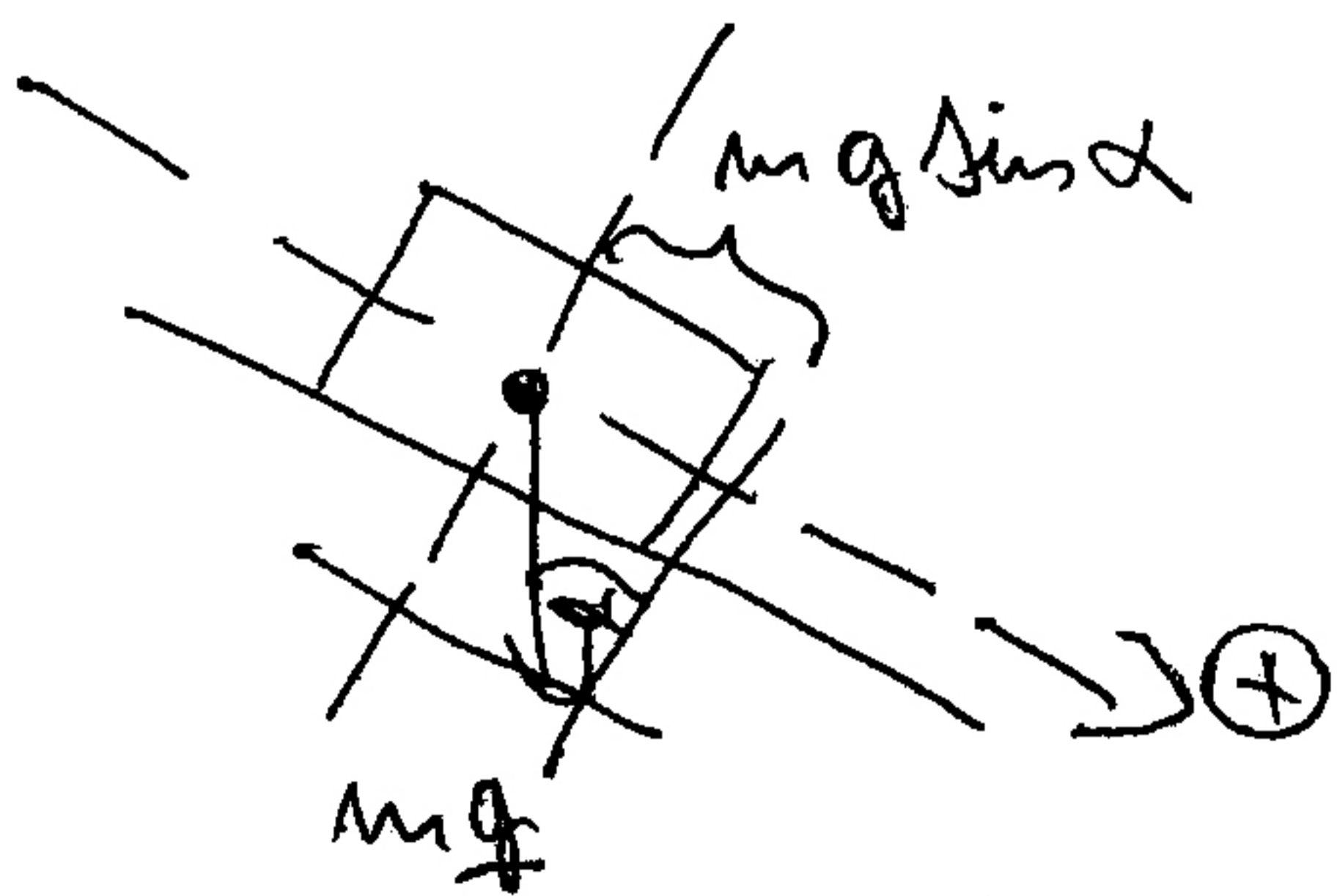
$$m \underline{a}' = -m \underline{a}_0 + m \underline{g} + \underline{N}$$

$$(m \underline{a}')_{||} = m \ddot{z}'$$

$$(-m \underline{a}_0)_{||} = -m a_0 \cos \alpha$$



$$(m \underline{g})_{||} = +m g \sin \alpha$$



$$(\underline{N})_{||} = 0$$

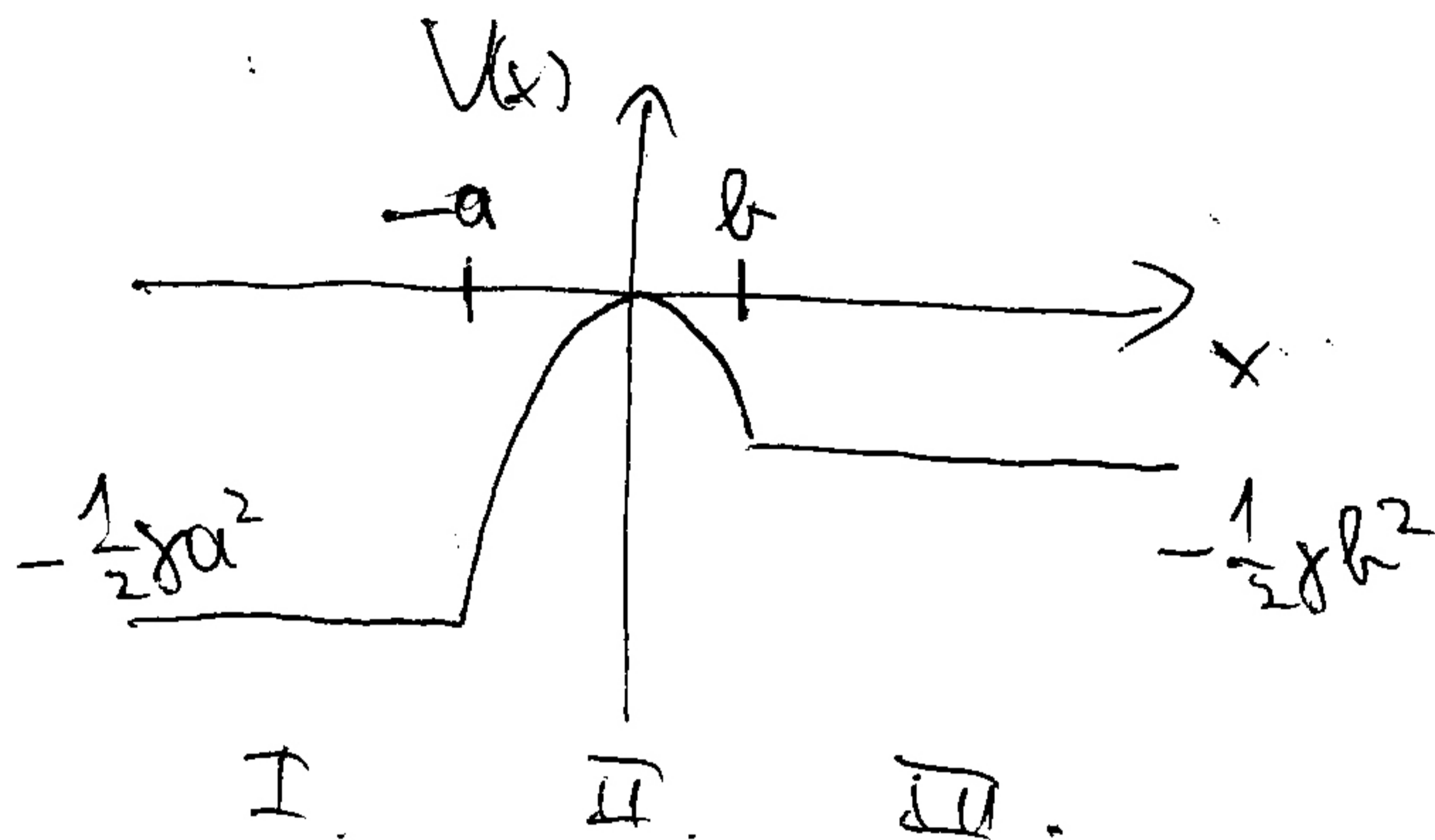
$$m \ddot{z}' = -m a_0 \cos \alpha + m g \sin \alpha$$

$$\ddot{z}' \stackrel{!}{=} 0$$

$$\rightarrow -m/a_0 \cos \alpha + m/g \sin \alpha \stackrel{!}{=} 0$$

$$|a| = a_0 \stackrel{!}{=} g \tan \alpha$$

$$2. \quad V(x) = \begin{cases} -\frac{1}{2}\gamma a^2, & \text{ha } x < -a \text{ I.} \\ -\frac{1}{2}\gamma x^2, & \text{ha } -a \leq x \leq b \text{ II.} \\ -\frac{1}{2}\gamma b^2, & \text{ha } x > b \text{ III.} \end{cases}$$



Egyensúlyi pont: $V'(x)|_{x=x^*} \stackrel{!}{=} 0$

I. $V'(x) = 0 \quad \forall x \in \text{I.}$

$\Rightarrow \forall x \in (-\infty, -a)$ egyensúlyi pont!

II. $V'(x) = -\gamma x$, $-\gamma x_{\text{II}}^* \stackrel{!}{=} 0 \rightarrow x_{\text{II}}^* = 0$

Itt ez az egy egyensúlyi pont van.

III. $V'(x) = 0 \quad \forall x \in \text{III.}$

$\Rightarrow \forall x \in (b, \infty)$ egyensúlyi pont.

Nyugalom: $\dot{x}|_{x=x^*} = 0, T \geq 0, V(x)|_{x=x^*} = E$

I. A teljes tartományon: $E = V(x) = -\frac{1}{2}\gamma a^2$
mellett nyugalomban lesz a tömegpont.

II. $E = V(x)|_{x=x_{II}^*} = -\frac{1}{2}\gamma x_{II}^{*2} = 0$

III. A teljes tartományon: $E = V(x) = -\frac{1}{2}\gamma b^2$
mellett nyugalomban lesz a tömegpont.

$a > b \Rightarrow -\frac{1}{2}\gamma a^2 < -\frac{1}{2}\gamma b^2 < 0$

Egy ilyen energiát tekintünk: ~~...~~

$$-\frac{1}{2}\gamma a^2 < E < -\frac{1}{2}\gamma b^2$$

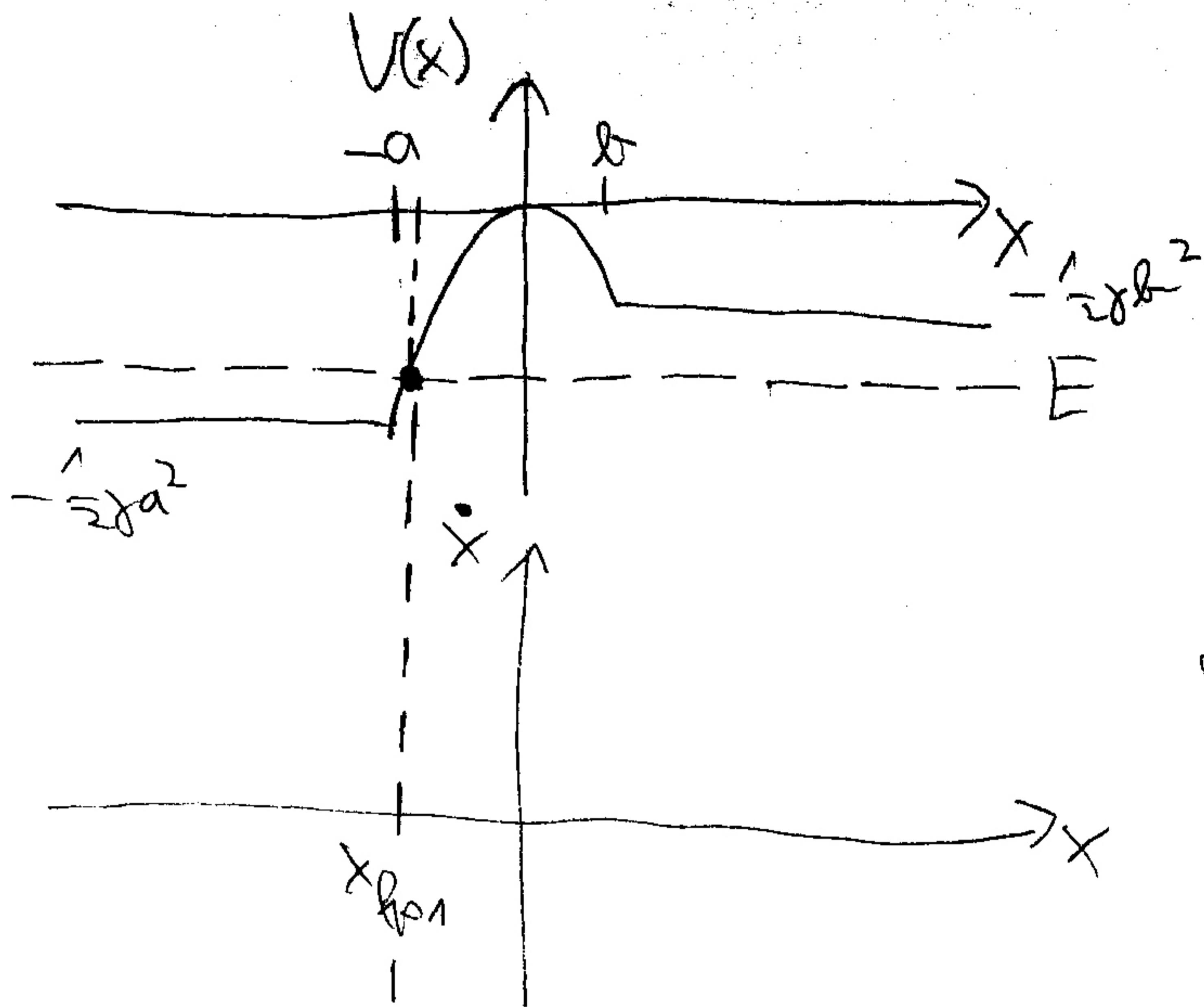
Mivel $E < V(x)|_{x=x_{II}^*}$ és $E < V(x)|_{x \in (b, \infty)}$

ezek az egyenletjű pontok nem tud áthaladni.

Mivel $E > V(x)|_{x \in (-\infty, -a)}$, az $x \in (-\infty, -a)$
egyenletjű pontok át tud haladni.

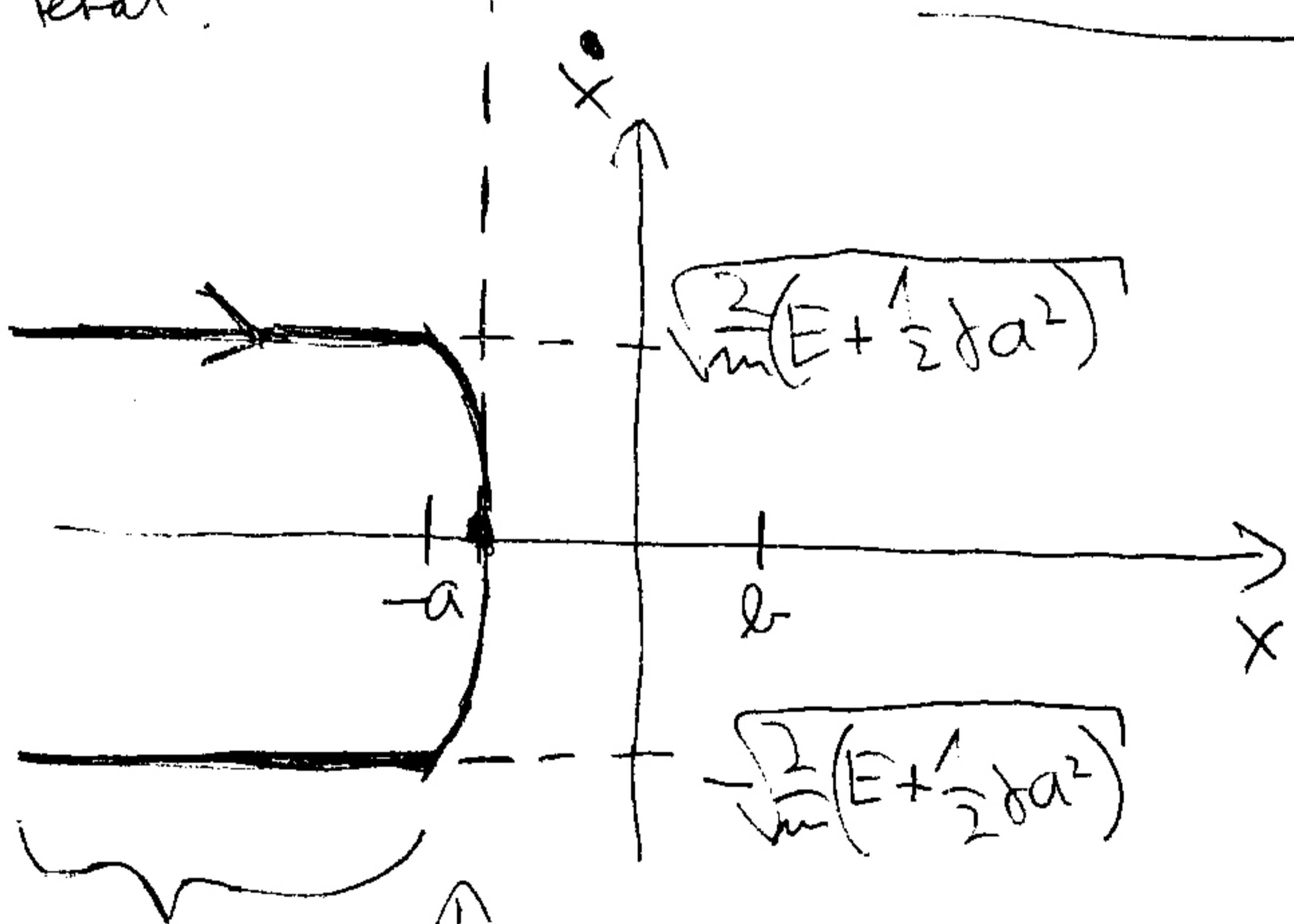
→ Utóbbiakban a sebesség:

$$\dot{x}|_{x \in (-\infty, a)} = \pm \sqrt{\frac{2}{m}(E - V(x)|_{x \in (-\infty, a)})} = \pm \sqrt{\frac{2}{m}(E + \frac{1}{2}\gamma a^2)}$$



Plathato, hogy
 ezzel a mozgással
 1 fordulópontja lesz.
 Ez a II. tartományba
 esik, azé:

Teljes:



(Itt egy kis körrel
 illeszkedik a trajektória.)

Ugyis látni, hogy
 itt minden pont lepusztul
 pont. Ezért balra át a tömegpont

$$\dot{x} = \pm \sqrt{\frac{2}{m} \left(E + \frac{1}{2} \gamma a^2 \right)}$$

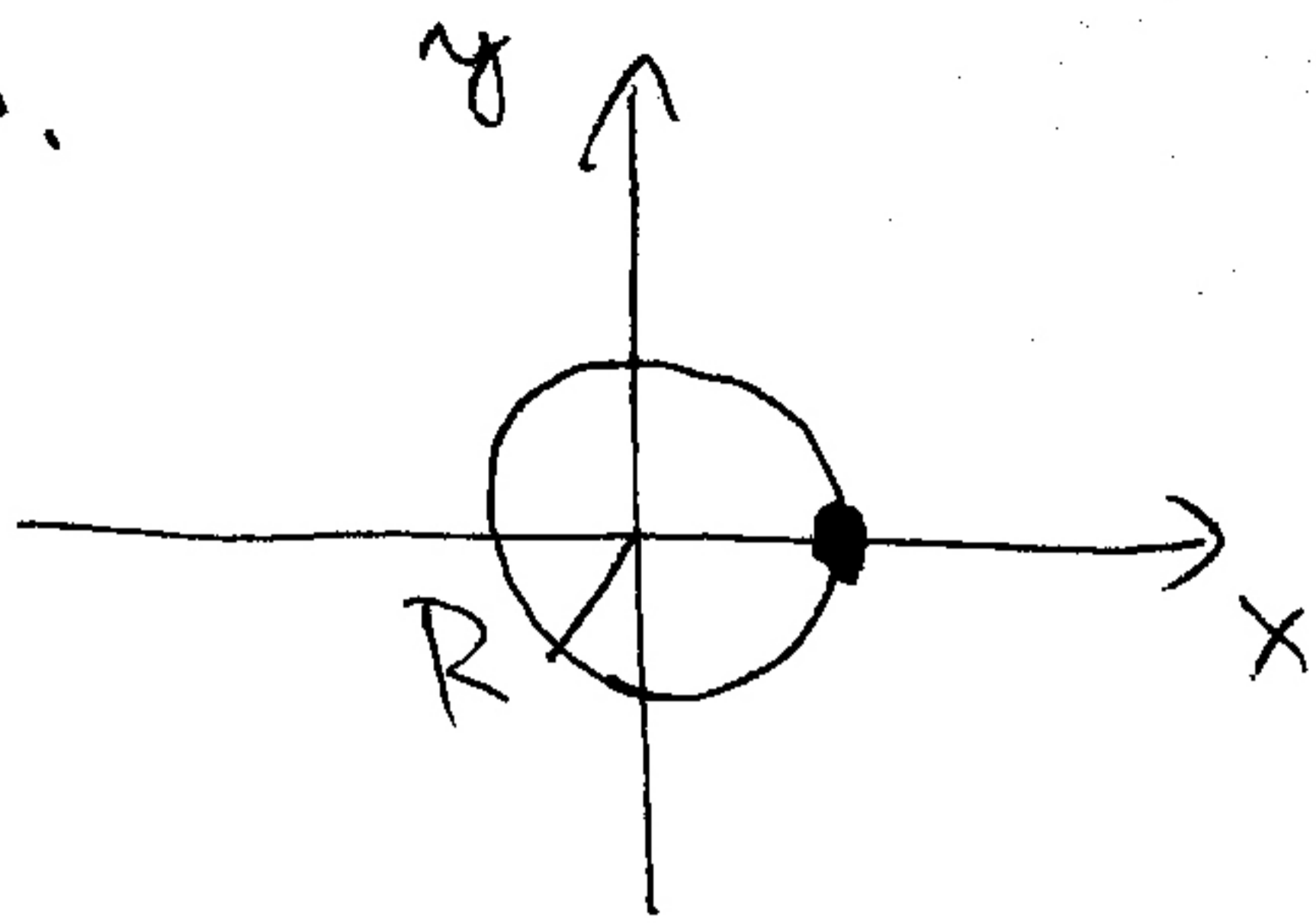
$$V(x_{fp}) = E$$

$$-\frac{1}{2} \gamma x_{fp}^2 = E$$

$$x_{fp1/2} = \pm \sqrt{\frac{2E}{\gamma}}$$

Mivel $x_{fp2} = \sqrt{\frac{2E}{\gamma}} > b$,
 ez a fordulópont nem
 valószínűleg meg. x_{fp1}
 viszont igen.

3.



$$v_{\pm r} = 0$$

$$\omega_{\pm} = \omega_0 \frac{R}{r}$$

$$t = \tau: v_{\text{rel}r} = 0, v_{\text{rel}\varphi} = 0,$$

$$x = R, y = 0 \Rightarrow r = R, \varphi = 0$$

Pérsygeben:

$$v_r(t) = ?, \varphi(t) = ?$$

$$v_{\text{rel}\varphi} = \cancel{\dots} 0$$

$$v_{\text{rel}r} = \tau^2 \frac{v_{\varphi}}{t^3}$$

$$v_r = v_{\pm r} + v_{\text{rel}r} = v_{\text{rel}r}$$

$$v_{\text{rel}r}(t) = v_{\text{rel}r}(\tau) + \int_{\tau}^t v_{\text{rel}r}(t') dt' = 0 + \int_{\tau}^t \tau^2 \frac{v_{\varphi}(t')}{t'^3} dt' \quad \begin{matrix} \uparrow \\ =? \end{matrix}$$

$$v_{\varphi} = v_{\pm\varphi} + v_{\text{rel}\varphi} = v_{\pm\varphi} = \omega_{\pm} \cdot r = \omega_0 \frac{R}{r} \cdot r = \omega_0 R$$

$$\rightarrow \boxed{v_r(t) = v_{\text{rel}r}(t) = \int_{\tau}^t \tau^2 \frac{\omega_0 R}{t'^3} dt' =}$$

$$= \tau^2 \omega_0 R \cdot \left(-\frac{1}{2}\right) \cdot \left(\frac{1}{t^2} - \frac{1}{\tau^2}\right) = \frac{1}{2} \omega_0 R - \frac{1}{2} \omega_0 R \frac{\tau^2}{t^2}$$

Es zeigt sich monoton wachsende fuggvény, így a limitét keressük.

$$\boxed{\lim_{t \rightarrow \infty} v_r(t) = \frac{1}{2} \omega_0 R}$$

ZH1P5t/6

$$\varphi(t) = \varphi(\tau) + \int_{\tau}^t \omega(t') dt'$$

$\searrow = ?$

$$\omega = \omega_f + \omega_{rel} = \omega_f + \frac{v_{rel}}{r} = \omega_f + \frac{0}{r} = \omega_f = \omega_0 + \frac{R}{r}$$

$$v(t) = v(\tau) + \int_{\tau}^t v_v(t') dt' = R + \int_{\tau}^t \left(\frac{1}{2} \omega_0 R - \frac{1}{2} \omega_0 R \frac{\tau^2}{t'^2} \right) dt' =$$

$v(t) = ?$

$$= R + \int_{\tau}^t \frac{1}{2} \omega_0 R (t' - \tau) dt' + \frac{1}{2} \omega_0 R \tau^2 \left(\frac{1}{t} - \frac{1}{\tau} \right) =$$

$$= R - \omega_0 R \tau + \frac{1}{2} \omega_0 R t + \frac{1}{2} \omega_0 R \frac{\tau^2}{t}$$

$$\rightarrow \varphi(t) = 0 + \int_{\tau}^t \omega_0 \cdot \frac{R}{v(t')} dt' =$$

Nenn lässt
R-t'ol.

↑

$$= \omega_0 \cdot \int_{\tau}^t \frac{R}{R - \omega_0 R \tau + \frac{1}{2} \omega_0 R t' + \frac{1}{2} \omega_0 R \frac{\tau^2}{t'}} dt' =$$

$$= \int_{\tau}^t \frac{t'}{\frac{1}{2} t'^2 + \left(\frac{1}{\omega_0} - \tau \right) t' + \frac{1}{2} \tau^2} dt' \quad (\text{Es ist integrierbar.})$$